# COL863: Quantum Computation and Information

Ragesh Jaiswal, CSE, IIT Delhi

# Quantum Computation: Quantum circuits

# Quantum Circuit Controlled operations

#### Theoerm

Suppose U is a unitary gate on a single qubit. Then there exist unitary operators A, B, C on a single qubit such that ABC = I and  $U = e^{i\alpha}AXBXC$ , where  $\alpha$  is some overall phase factor.

#### Question

For a single qubit U, can we implement Controlled-U gate using only CNOT and single-qubit gates? Yes

### Construction sketch

The construction follows from the following circuit equivalences.



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For a single qubit U, can we implement Controlled-U gate with n control qubits using only CNOT and single-qubit gates?

# Quantum Circuit Controlled operations

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For a single qubit U, can we implement Controlled-U gate with n control qubits using only CNOT and single-qubit gates? Yes using ancilla qubits

#### Construction sketch



• A few other gates and circuit identities:



Figure: NOT gate applied to the target qubit conditional on the control qubit being 0.



## Principle of deferred measurements

Measurements can always be moved from an intermediate stage of a quantum circuit to the end of the circuit; if the measurement results are used at any stage of the circuit, then the clasically controlled operations can be replaced by conditional quantum operations.



# Quantum Circuit Measurements

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#### Principle of implicit measurement

Without loss of generality, any unterminated quantum wires (qubits which are not measured) at the end of a quantum circuit may be assumed to be measured.

• Exercise: Suppose  $\rho$  is the density matrix describing a two qubit system. Suppose we perform a projective measurement in the computational basis of the second qubit. Let  $P_0 = I \otimes |0\rangle \langle 0|$  and  $P_1 = I \otimes |1\rangle \langle 1|$  be the projectors onto the  $|0\rangle$  and  $|1\rangle$  states of the second qubit, respectively. Let  $\rho'$  be the density matrix which would be assigned to the system after the measurement by an observer who did not learn the measurement result. Show that

$$\rho' = P_0 \rho P_0 + P_1 \rho P_1.$$

Also show that the reduced density matrix for the first qubit is not affected by the measurement, that is,  $tr_2(\rho) = tr_2(\rho')$ .

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## Principle of implicit measurement

Without loss of generality, any unterminated quantum wires (qubits which are not measured) at the end of a quantum circuit may be assumed to be measured.

• Exercise: Show that measurement commutes with control.



• A set of gates is said to be universal for quantum computation if any unitary operation may be **approximated** to arbitrary accuracy by a quantum circuit involving only those gates.

### Claim

Any unitary operation can be approximated to arbitrary accuracy using Hadamard, phase, CNOT, and  $\pi/8$  gates.

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### Proof sketch

- <u>Claim 1</u>: A single qubit operation may be approximated to arbitrary accuracy using the Hadamard, phase, and π/8 gates.
- <u>Claim 2</u>: An arbitrary unitary operator may be expressed **exactly** using single qubit and CNOT gates.
  - <u>Claim 2.1</u>: An arbitrary unitary operator may be expressed **exactly** as a product of unitary operators that each acts non-trivially only on a subspace spanned by two computational basis states (such gates are called two-level gates).
  - Claim 2.2: An arbitrary two-level unitary operator may be expressed exactly using using single qubit and CNOT gates.
- What about efficiency?
  - Upper-bound: Any unitary can be approximated using exponentially many gates.
  - Lower-bound: There exists a unitary operation that which require exponentially many gates to approximate.

## Claim 2.1

An arbitrary unitary operator may be expressed **exactly** as a product of unitary operators that each acts non-trivially only on a subspace spanned by two computational basis states.

## Proof sketch

• The main idea can be understood using a  $3 \times 3$  unitary matrix:

$$U = \begin{bmatrix} a & d & g \\ b & e & h \\ c & f & j \end{bmatrix}$$

• We will find two-level unitary matrices  $U_1, U_2, U_3$  such that

$$U_3U_2U_1U = I$$
 and  $U = U_1^{\dagger}U_2^{\dagger}U_3^{\dagger}$ 

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### • Exercise

- Show that any  $d \times d$  unitary matrix can be written in terms of d(d-1)/2 two-level matrices.
- There exists a  $d \times d$  unitary matrix U which cannot be decomposed as a product of fewer than d-1 two-level unitary matrices.

#### Claim 2

An arbitrary unitary operator may be expressed **exactly** using single qubit and CNOT gates.

- <u>Claim 2.1</u>: An arbitrary unitary operator may be expressed exactly as a product of unitary operators that each acts non-trivially only on a subspace spanned by two computational basis states.
- <u>Claim 2.2</u>: An arbitrary two-level unitary operator may be expressed exactly using using single qubit and CNOT gates.

#### Proof sketch

- Let U be a two-level unitary matrix on a n-qubit quantum computer.
- Let U act non-trivially on the space spanned by the computational basis states |s⟩ and |t⟩, where s = s<sub>1</sub>,..., s<sub>n</sub> and t = t<sub>1</sub>,..., t<sub>n</sub> are n-bit binary strings.
- Let *Ũ* be the non-trivial 2 × 2 submatrix of *U*. Note that we can think *Ũ* to be a unitary operator on a single qubit.
- We will use the gray-code connecting s and t which is a sequence of n-bit strings staring with s and ending with t such that the subsequent strings in the sequence differ only on one bit.
- Example: *s* = 101001, *t* = 110011.

 $g_1 = 101001; g_2 = 101011; g_3 = 100011; g_4 = 110011$ 

- Main idea:
  - . We will design a sequence of swaps
    - $|g_1\rangle \rightarrow |g_{m-1}\rangle, |g_2\rangle \rightarrow |g_1\rangle, |g_3\rangle \rightarrow |g_2\rangle, ..., |g_{m-1}\rangle \rightarrow |g_{m-2}\rangle.$
  - . We will apply  $\tilde{U}$  to the qubit that differs in  $g_{m-1}$  and  $g_m$ .
  - $_{\diamond}$  Swap  $|g_{m-1}\rangle$  with  $|g_{m-2}\rangle,\,|g_{m-2}\rangle$  with  $|g_{m-3}\rangle$  and so on.

## Claim 2.2

An arbitrary two-level unitary operator may be expressed exactly using using single qubit and CNOT gates.

### Example construction

• Let the two-level transformation be:

$$U = \begin{bmatrix} a & 0 & 0 & 0 & 0 & 0 & 0 & c \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ b & 0 & 0 & 0 & 0 & 0 & 0 & d \end{bmatrix}$$

• The gray code connecting  $|000\rangle$  and  $|111\rangle$ :  $|000\rangle \rightarrow |001\rangle \rightarrow |011\rangle \rightarrow |111\rangle$ .

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- Exercise
  - For an arbitrary unitary operator on an *n*-qubit system, how many CNOT and single qubit gate will be required in the entire construction?

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#### Example construction

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U

Construction:



- Exercise
  - For an arbitrary unitary operator on an *n*-qubit system, how many CNOT and single qubit gate will be required in the entire construction?  $O(n^24^n)$  gates.

Any unitary operation can be approximated to arbitrary accuracy using Hadamard, phase, CNOT, and  $\pi/8$  gates.

### Proof sketch

- <u>Claim 1</u>: A single qubit operation may be approximated to arbitrary accuracy using the Hadamard, phase, and π/8 gates.
- <u>Claim 2</u>: An arbitrary unitary operator may be expressed **exactly** using single qubit and CNOT gates.
  - <u>Claim 2.1</u>: An arbitrary unitary operator may be expressed **exactly** as a product of unitary operators that each acts non-trivially only on a subspace spanned by two computational basis states (such gates are called two-level gates).
  - <u>Claim 2.2</u>: An arbitrary two-level unitary operator may be expressed exactly using using single qubit and CNOT gates.
- A discrete set of gates cannot be used to implement an arbitrary unitary operation.
- However, it may be possible to approximate any unitary gate using a discrete set of gates.

A single qubit operation may be **approximated** to arbitrary accuracy using the Hadamard, phase, and  $\pi/8$  gates.

- We first need to define a notion of approximating a unitary operation.
- Let U and V be unitary operators on the same state space.
  - *U* denotes the target unitary operator that we would like to implement.
  - V is the operator that is actually implemented.
- The error (w.r.t. implementing V instead of U) is defined as

$$E(U,V)\equiv \max_{\ket{\psi}} \ket{\ket{U-V}\ket{\psi}}$$

• Question: Why is the above a reasonable notion of error when implementing V instead of U?

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• The error (w.r.t. implementing V instead of U) is defined as

$$E(U, V) \equiv \max_{\ket{\psi}} ||(U - V) \ket{\psi}||$$

## Claim 1.1

Suppose we wish to implement a quantum circuit with *m* gates  $U_1, ..., U_m$ . However, we can only implement  $V_1, ..., V_m$ . The difference in probabilities of a measurement outcome will be at most a tolerance  $\Delta > 0$  given that  $\forall j, E(U_j, V_j) \leq \frac{\Delta}{2m}$ .

### Claim 1

A single qubit operation may be **approximated** to arbitrary accuracy using the Hadamard, phase, and  $\pi/8$  gates.

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### Proof sketch

- <u>Claim 1.1.1</u>: For any POVM element *M* let *P<sub>U</sub>* and *P<sub>V</sub>* denote the probabilities for measuring this element when *U* and *V* are used respectively. Then |*P<sub>U</sub>* − *P<sub>V</sub>*| ≤ 2 · *E*(*U*, *V*).
- <u>Claim 1.1.2</u>:  $E(U_m U_{m-1} ... U_1, V_m V_{m-1} ... V_1) \le \sum_{j=1}^m E(U_j, V_j).$

# End

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