

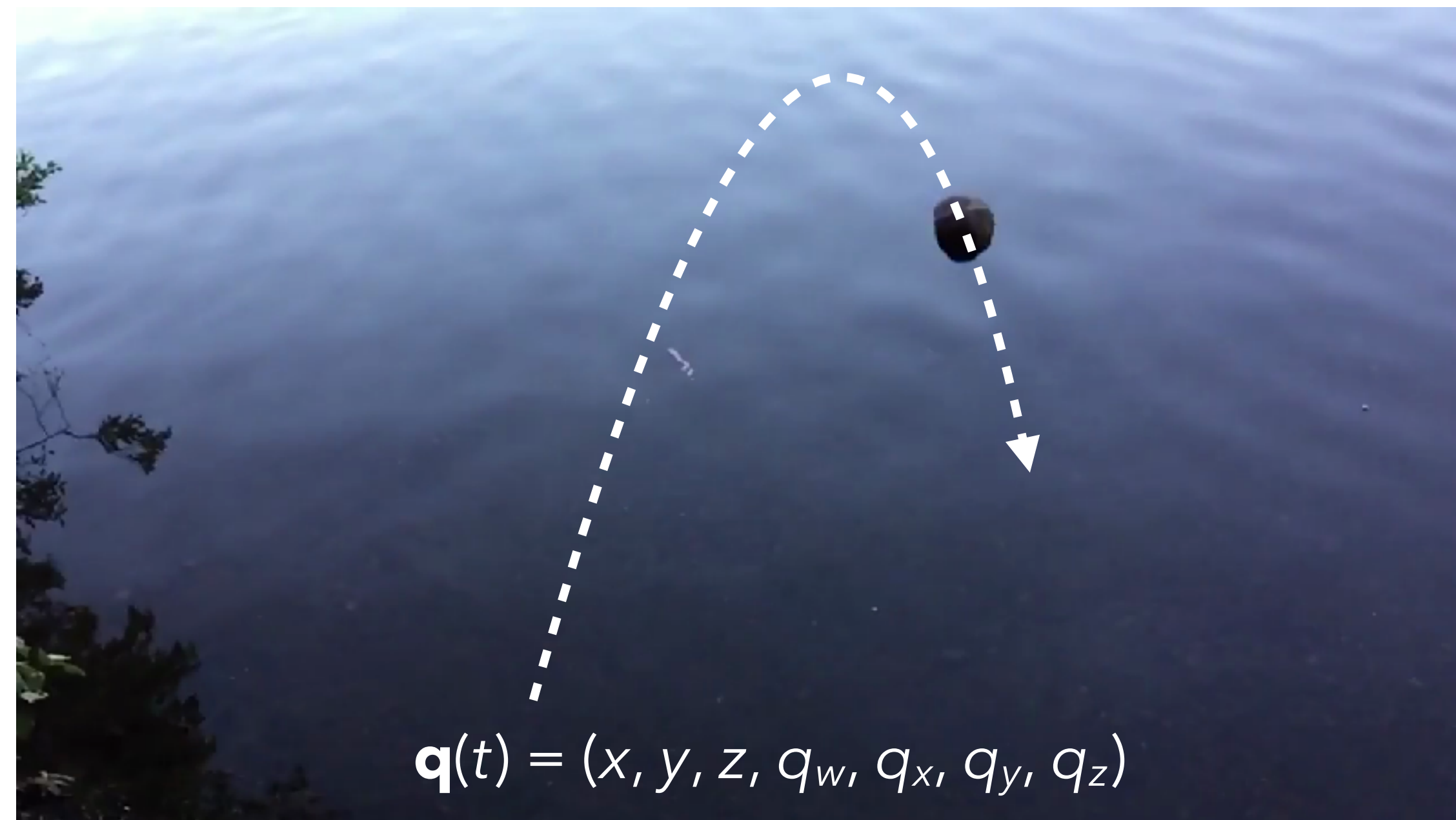


COL781: Computer Graphics

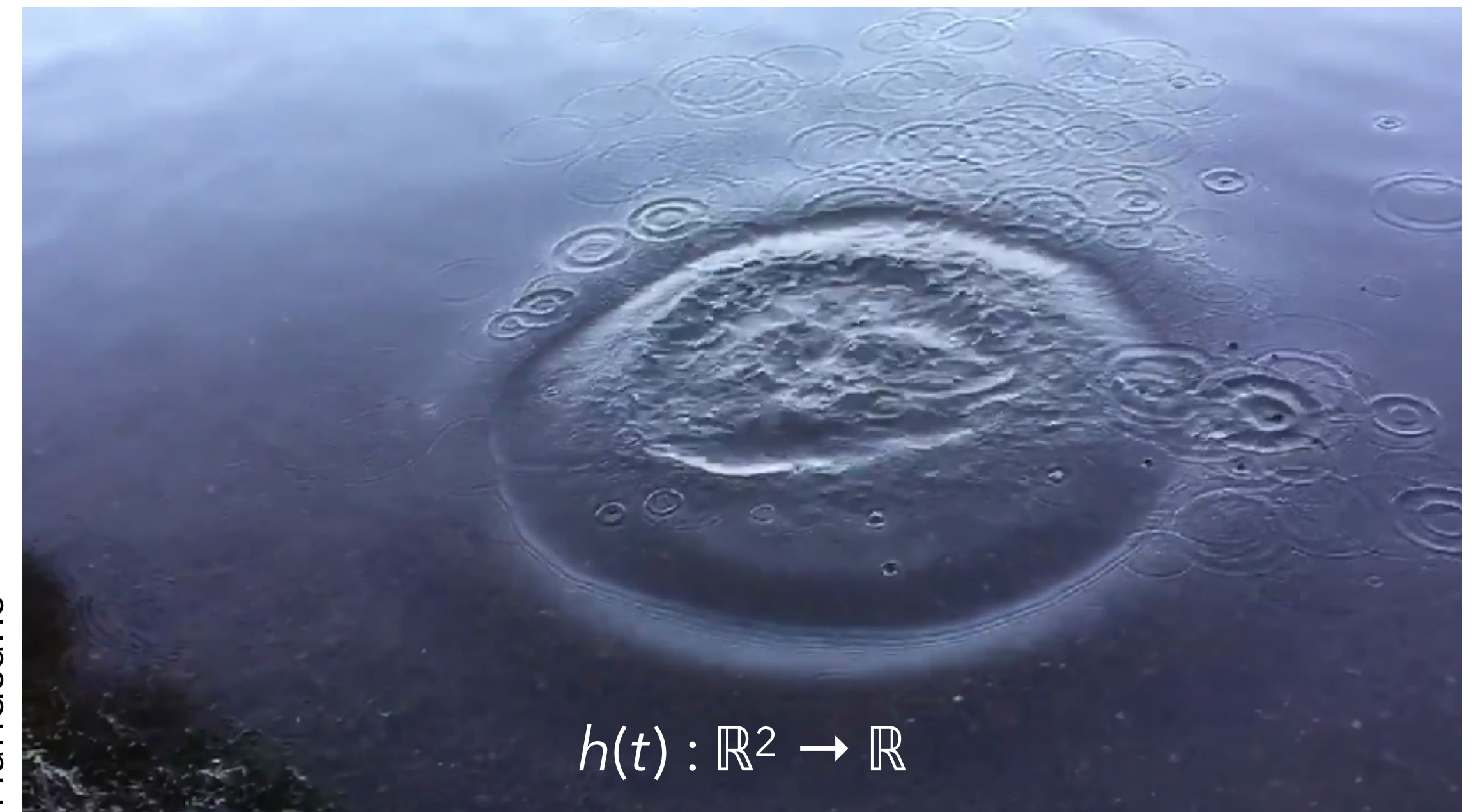
35. Continuum Models

So far, we know how to simulate discrete systems of particles and rigid bodies.

But lots of things in real life are not discrete:



Handsans





https://www.youtube.com/watch?v=RRPP73QM_4k

Liquids



<https://www.youtube.com/watch?v=jVxYuPeqOPI>

Smoke



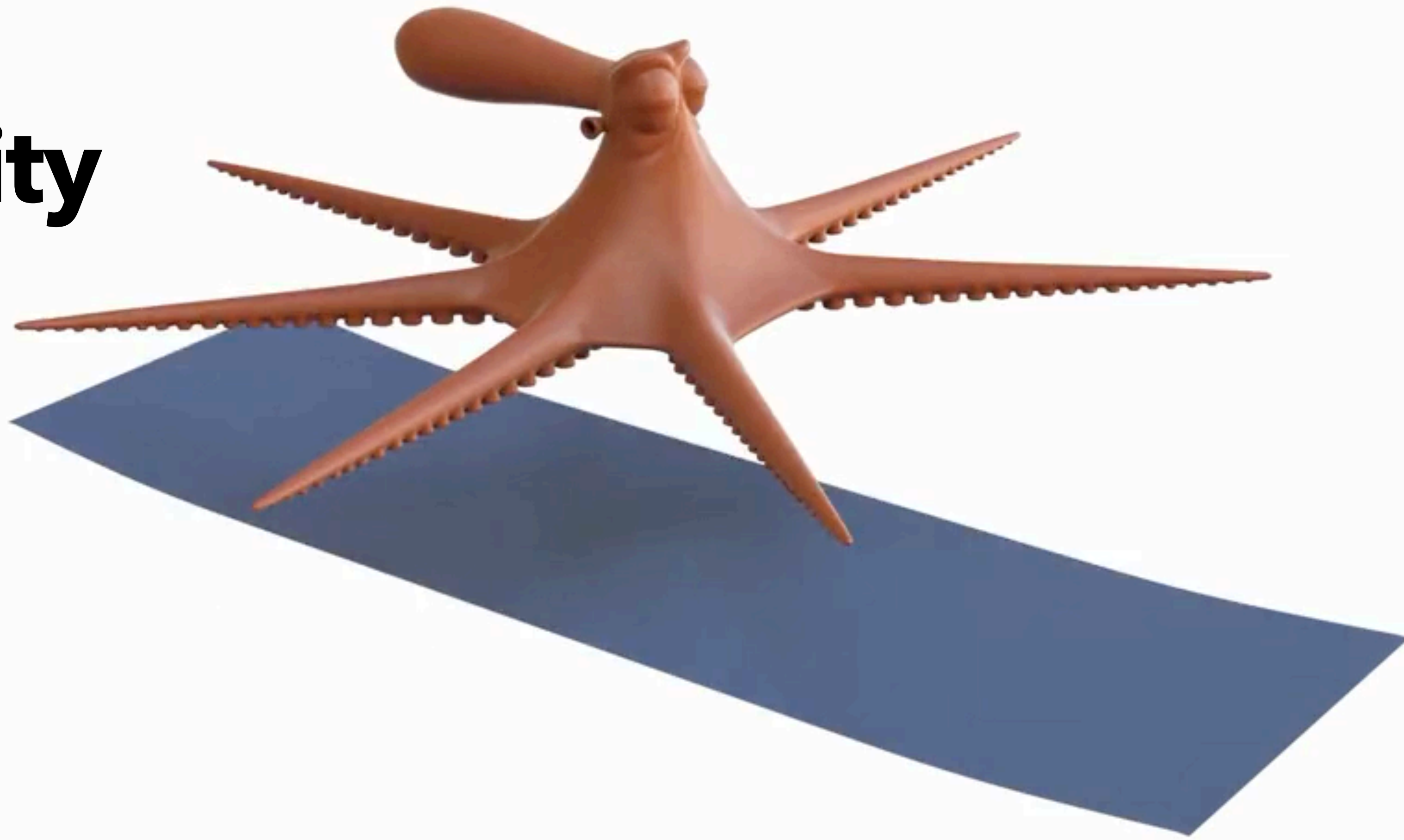
[https://research.nvidia.com/publication/
2008-07_low-viscosity-flow-simulations-animation](https://research.nvidia.com/publication/2008-07_low-viscosity-flow-simulations-animation)

Cloth



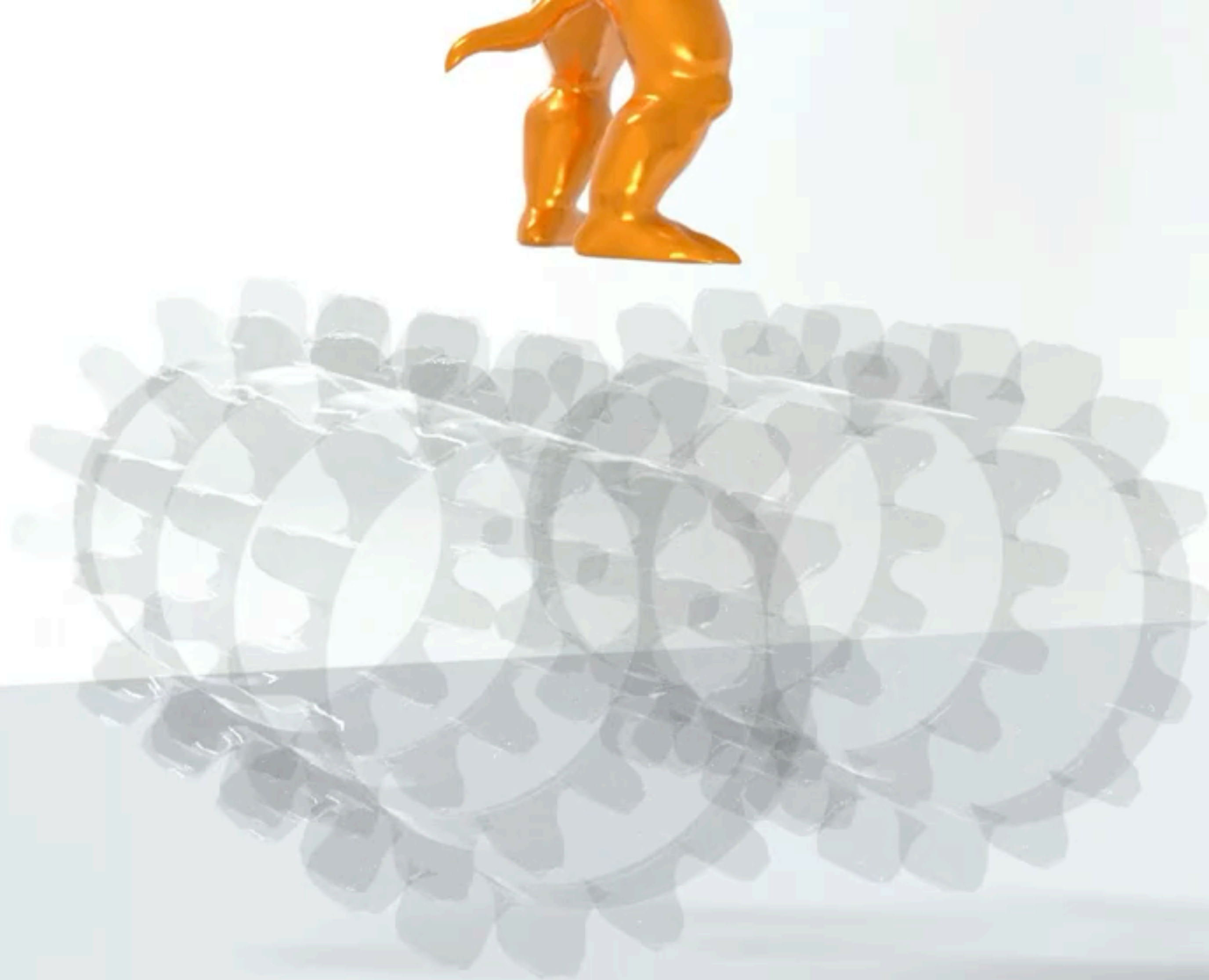
<https://www.youtube.com/watch?v=TH5g8TuKlkk>

Elasticity



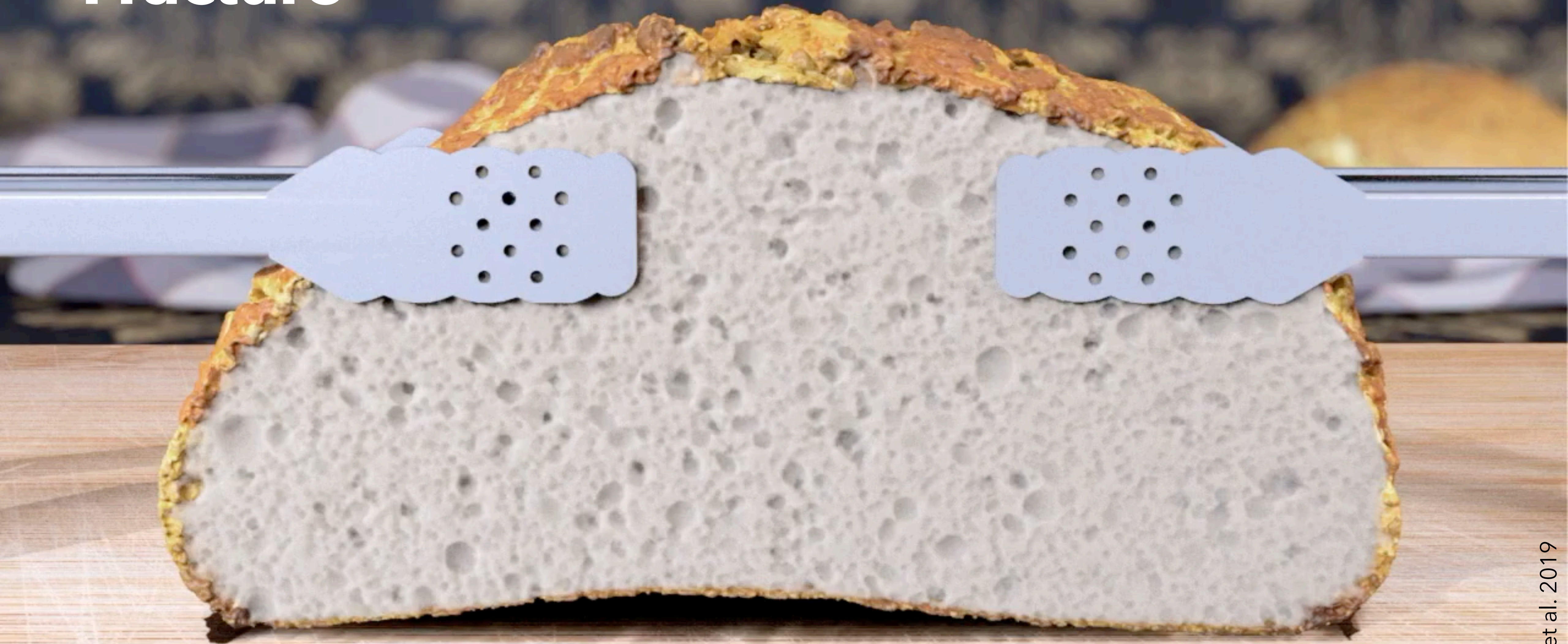
<https://vimeo.com/333798247>

Plasticity



<https://www.youtube.com/watch?v=YvvoSu8NK3A>

Fracture



<https://www.youtube.com/watch?v=INri-x2nK7o>

Snow



<https://vimeo.com/160322962>

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Stomakhin et al. 2013

Mass-spring chain \rightarrow continuum rod

- Positions: $(\dots, \mathbf{x}_i, \dots)$ where $i \in \{1, 2, \dots, n\}$
- Mass of particle: m_i
- Stretch of spring: $\|\mathbf{x}_{i+1} - \mathbf{x}_i\|/\ell_0$
- Force on particle: $\mathbf{f}_i - \mathbf{f}_{i-1}$
- Positions: $\mathbf{x}(s)$ where $s \in [0, L]$
- Mass of differential segment: $dm = \rho ds$
- Stretch of segment: $\|\mathbf{x}(s+ds) - \mathbf{x}(s)\|/ds = \|\partial\mathbf{x}/\partial s\|$
- Differential force on segment: $\mathbf{f}(s+ds) - \mathbf{f}(s) = (\partial\mathbf{f}/\partial s) ds$

Equations of motion: ODE

$$d^2\mathbf{x}_i/dt^2 = m_i^{-1} \mathbf{f}(\mathbf{x}_i, \mathbf{x}_{i-1}, \mathbf{x}_{i+1}, \dots)$$

Equations of motion: partial differential eq.

$$\partial^2\mathbf{x}/\partial t^2 = \rho^{-1} \mathbf{f}(\mathbf{x}, \partial\mathbf{x}/\partial s, \partial^2\mathbf{x}/\partial s^2, \dots)$$

$$\frac{\partial^2 \mathbf{x}}{\partial t^2} = \rho^{-1} \mathbf{f} \left(\mathbf{x}, \frac{\partial \mathbf{x}}{\partial s}, \frac{\partial^2 \mathbf{x}}{\partial s^2}, \dots \right)$$

Our goal in this course:

- Not to fully understand the physics (except enough to gain intuition)
- Not to fully understand the mathematics (except enough to gain intuition)
- Understand how to compute **numerical solutions** to such equations!

We'll start with some simpler examples...

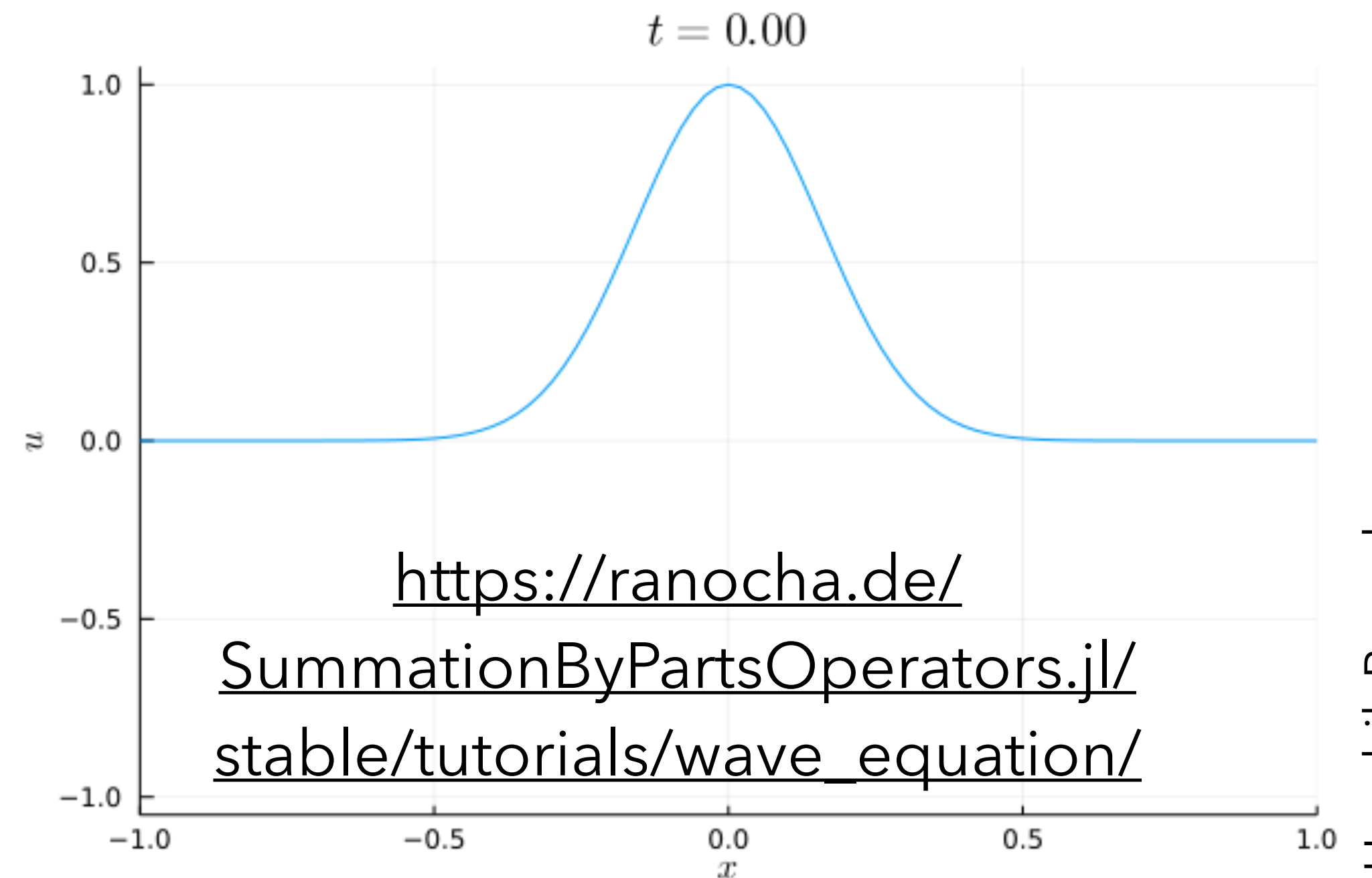
The wave equation

$$\frac{\partial^2 u}{\partial t^2} = \frac{\partial^2 u}{\partial x^2}$$

Similar to a harmonic oscillator $d^2u/dt^2 = -ku$:
leads to oscillations

Right-hand side d^2u/dx^2 is curvature of graph $u(x)$

- Restoring force tries to straighten curves, flatten extrema



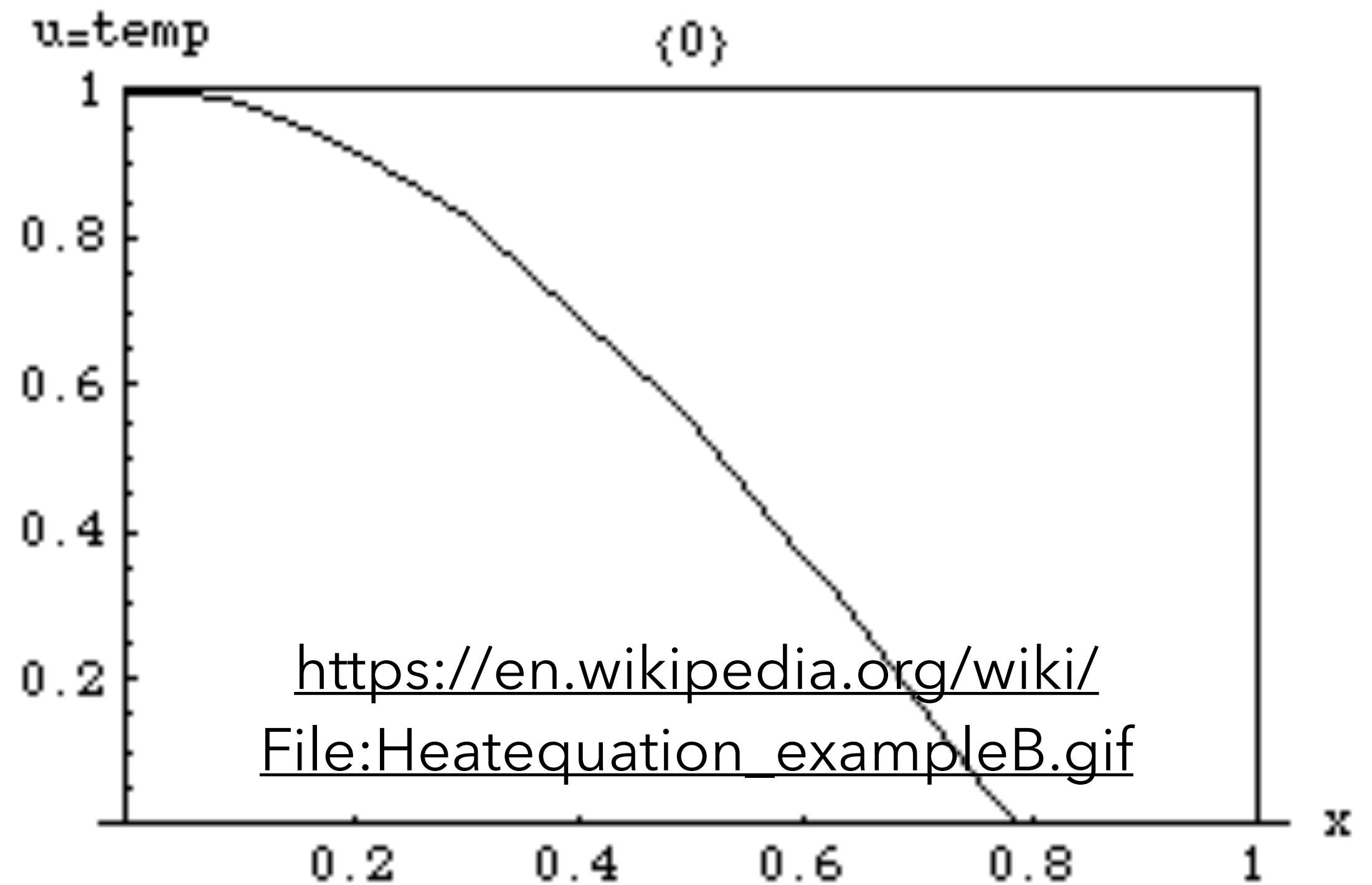
The heat equation

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}$$

- Second derivative \rightarrow oscillations / waves
- First derivative \rightarrow decay / smoothing

So, here is our problem:

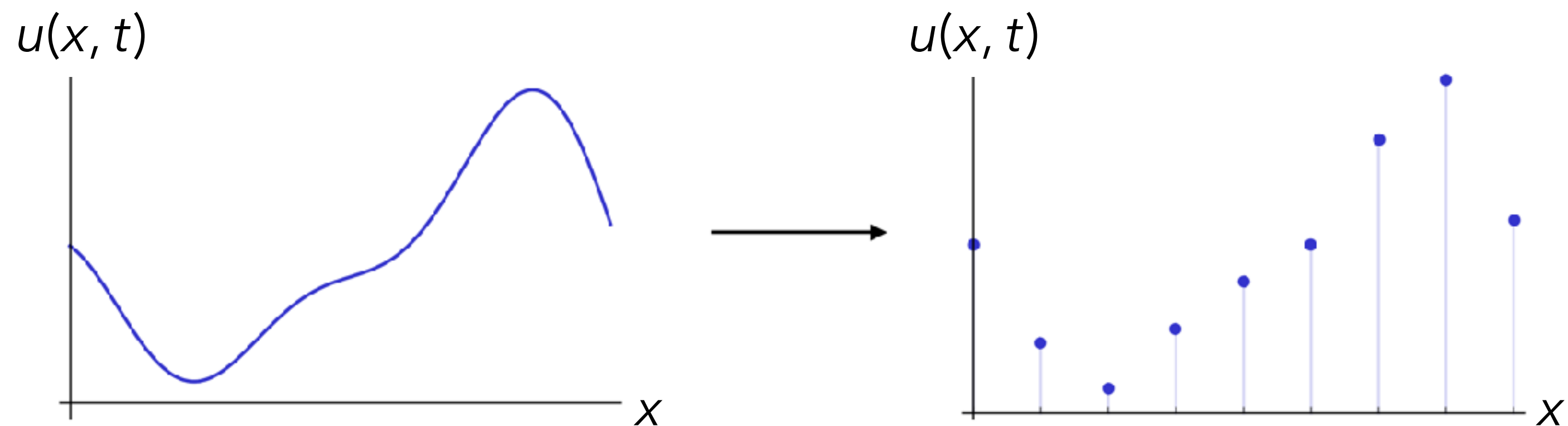
Given the partial differential equation
and the initial conditions $u(x, t=0)$,
find the spatial distribution $u(x, t)$ at all future times t



At any future time t , how to represent the function $u(x, t)$?

- Even if $u(x, 0)$ has an analytical form, $u(x, t)$ probably doesn't for $t > 0$

Simplest representation: store samples $u_i(t) = u(x_i, t)$ at various x_i



Now we need to solve $\frac{du_i}{dt} = \frac{\partial^2 u}{\partial x^2} \dots$

How to estimate spatial derivatives from discrete samples?

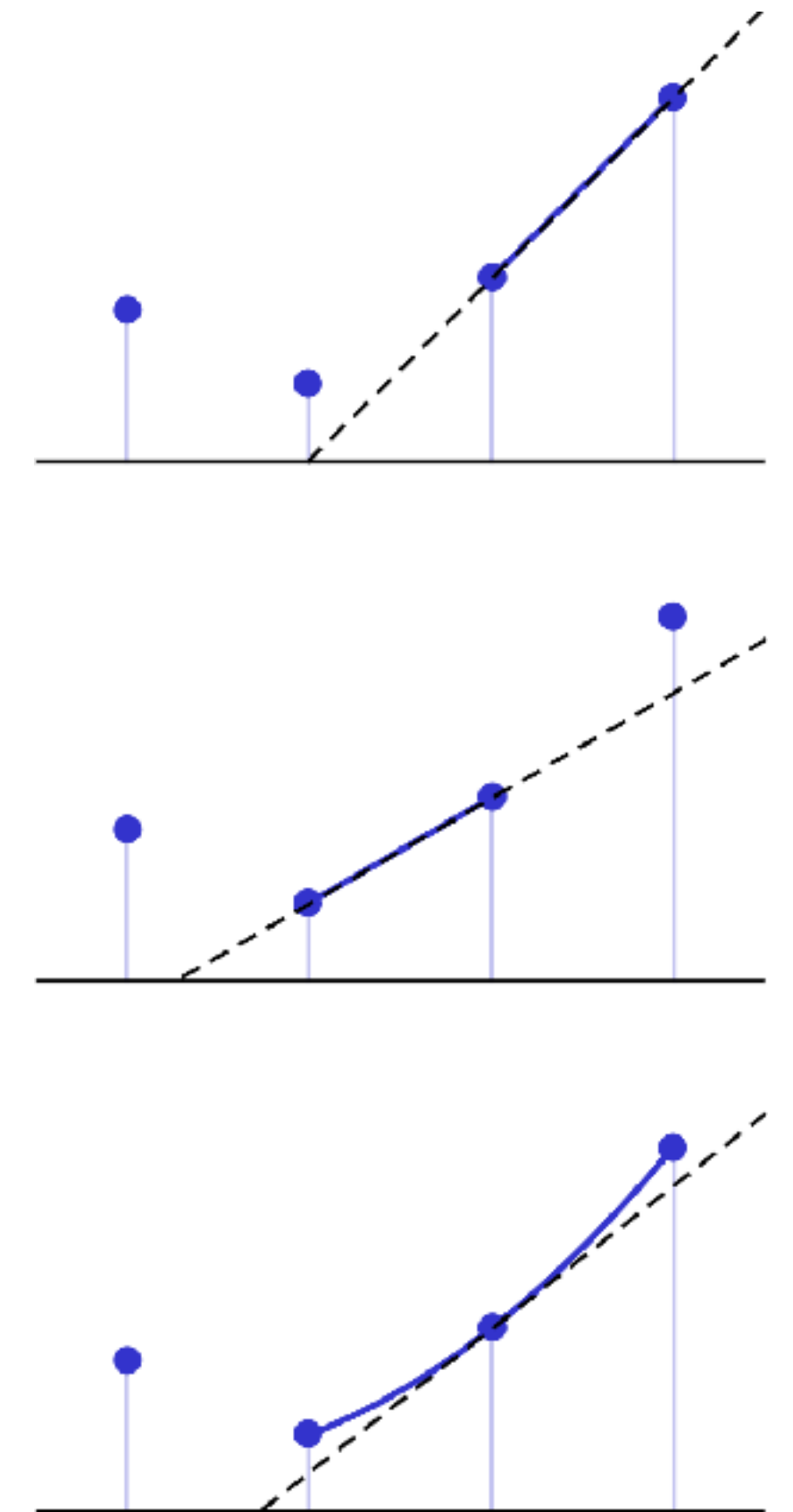
General strategy: reconstruct a sufficiently smooth function, then differentiate it

$$\frac{du}{dx} \approx \frac{u_{i+1} - u_i}{\Delta x} \quad (\text{forward difference})$$

$$\approx \frac{u_i - u_{i-1}}{\Delta x} \quad (\text{backward difference})$$

$$\approx \frac{u_{i+1} - u_{i-1}}{2\Delta x} \quad (\text{centered difference})$$

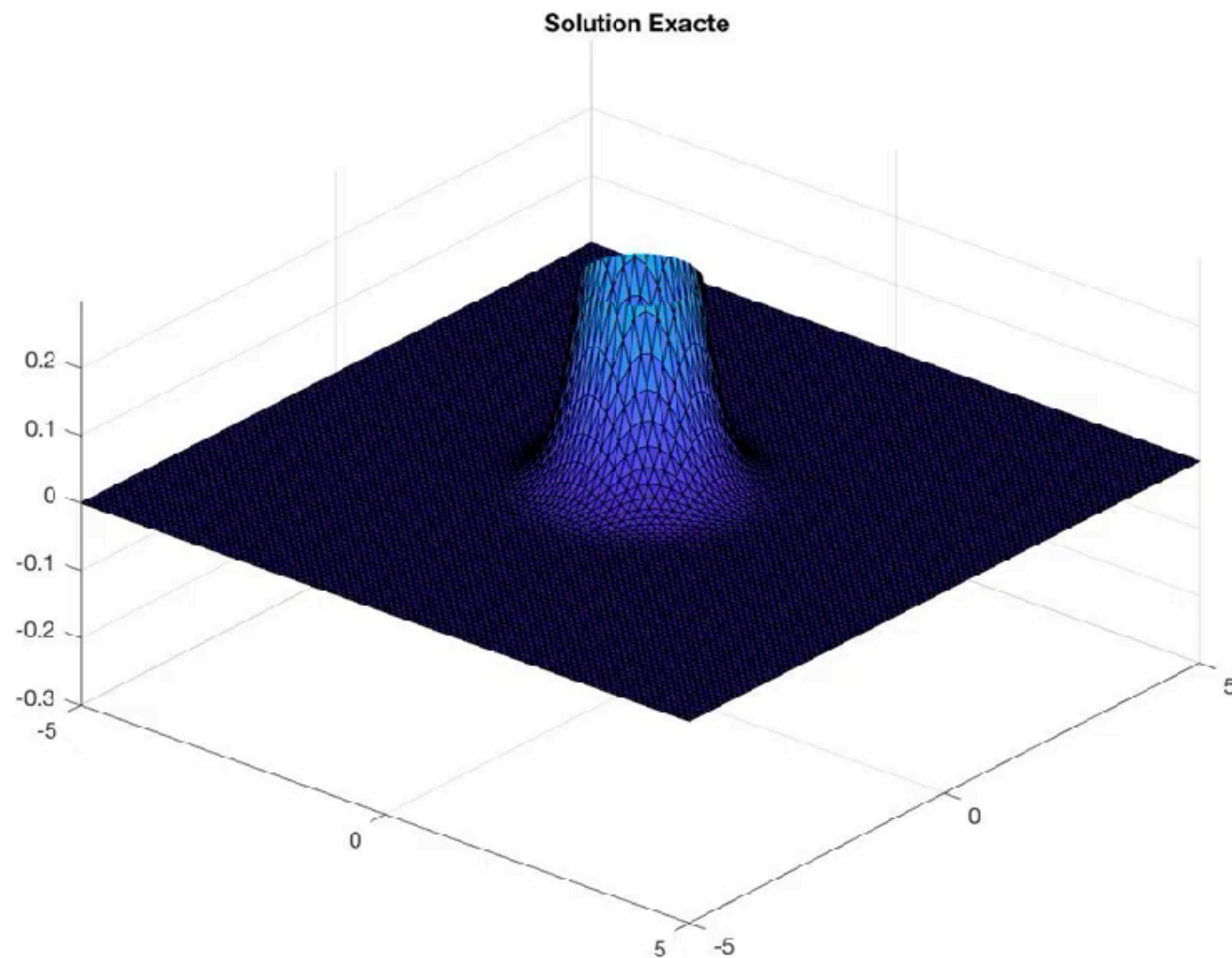
$$\frac{d^2u}{dx^2} \approx \frac{u_{i-1} - 2u_i + u_{i+1}}{\Delta x^2}$$



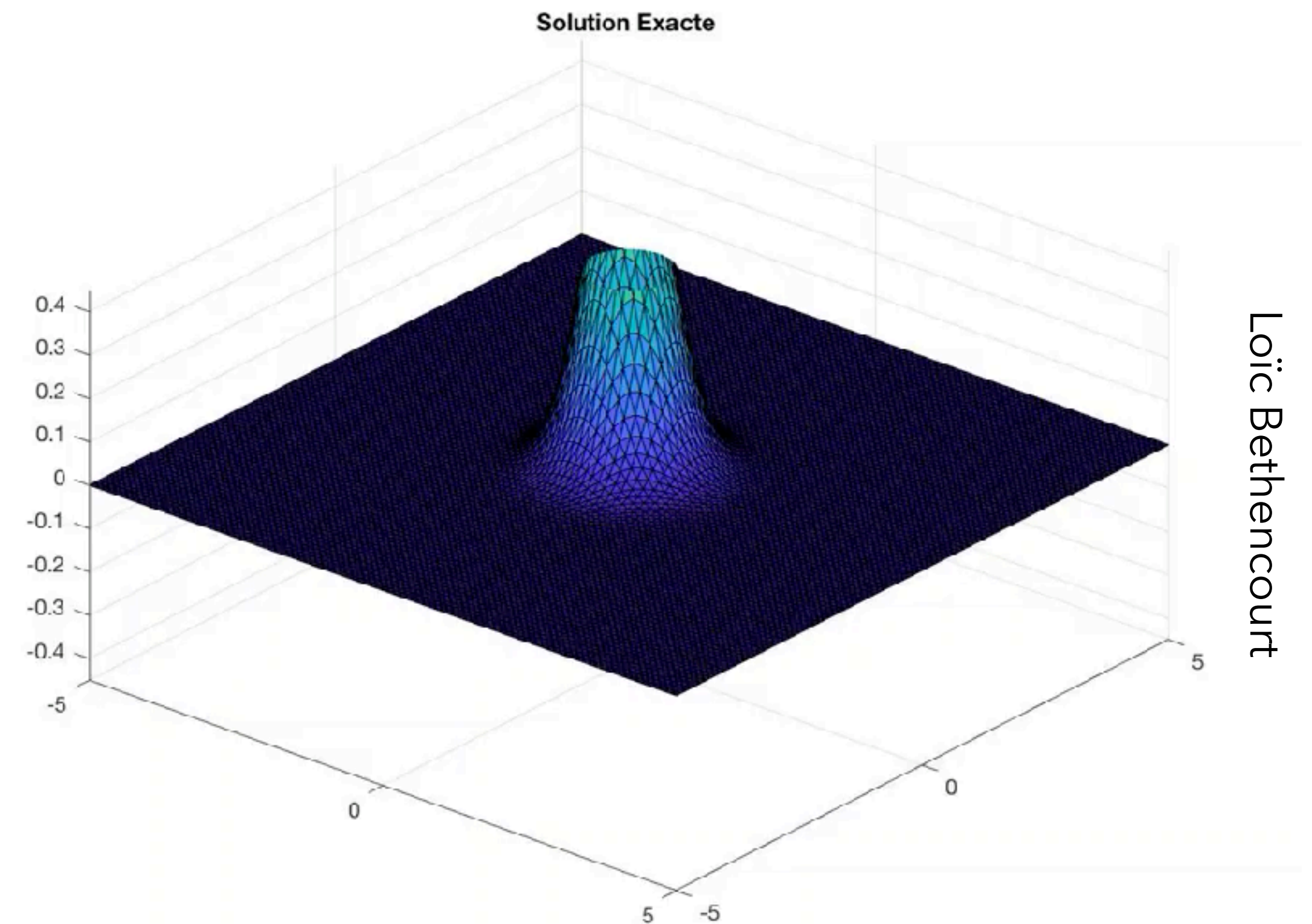
Also gives intuition: $d^2u/dx^2 \propto (\text{average of neighbours}) - u_i$

Boundary conditions

Dirichlet boundaries: $u = \text{fixed}$



Neumann boundaries: $\mathbf{n} \cdot \nabla u = \text{fixed}$



<https://www.youtube.com/watch?v=-chMgHvZxH0>

<https://www.youtube.com/watch?v=1hsj10dOgt0>

Anyway, we can't do things like $\frac{u_{i-1} - 2u_i + u_{i+1}}{\Delta x^2}$ if $i = 0$ or N .

- Dirichlet boundary: $u = \text{user-specified } f$

Easy: Just fix $u_0 = f$

- Neumann boundary: $du/dx = \text{user-specified } g$

Create a "ghost node" u_{-1} so that $\frac{u_0 - u_{-1}}{\Delta x} = g$, then plug in

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}$$

↓

$$\frac{du_i}{dt} = \frac{\partial^2 u}{\partial x^2}$$

↓

$$\frac{du_i}{dt} = \frac{u_{i-1} - 2u_i + u_{i+1}}{\Delta x^2}$$

This is now an ordinary differential equation! And we know how to solve those.

- Choose a time integration scheme, solve the equations, ...

PDEs in higher dimensions

In 1D, unknown function is $u = u(x, t)$, equation is something like $\frac{\partial^2 u}{\partial t^2} = f\left(u, \frac{\partial u}{\partial x}, \frac{\partial^2 u}{\partial x^2}, \dots\right)$

In n D, it's $u = u(x, y, \dots, t)$

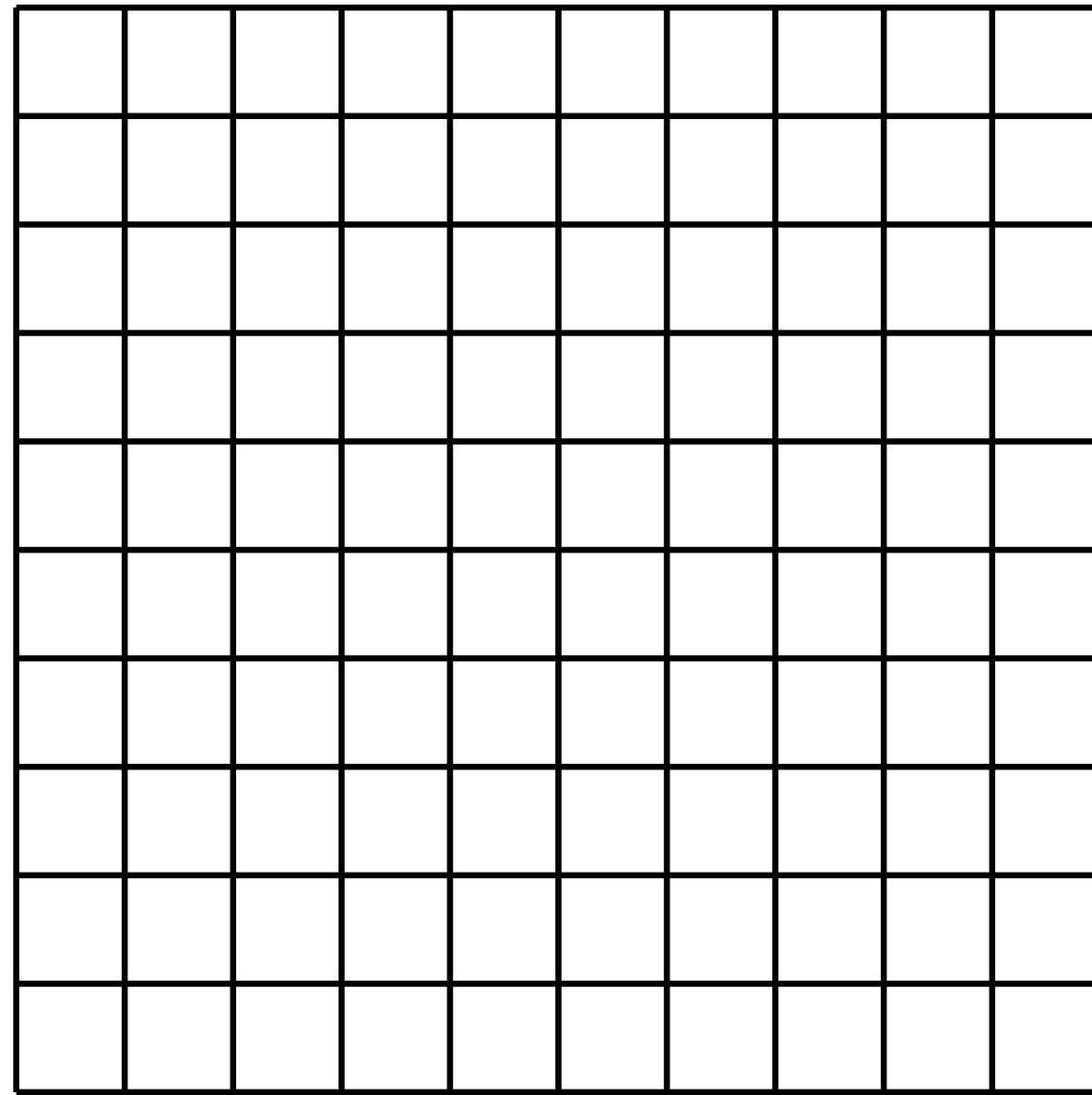
$$\frac{\partial^2 u}{\partial t^2} = f\left(u, \frac{\partial u}{\partial x}, \frac{\partial u}{\partial y}, \dots, \frac{\partial^2 u}{\partial x^2}, \frac{\partial^2 u}{\partial y^2}, \frac{\partial^2 u}{\partial x \partial y}, \dots\right)$$

In particular, right-hand side of heat equation / wave equation becomes the **Laplacian**

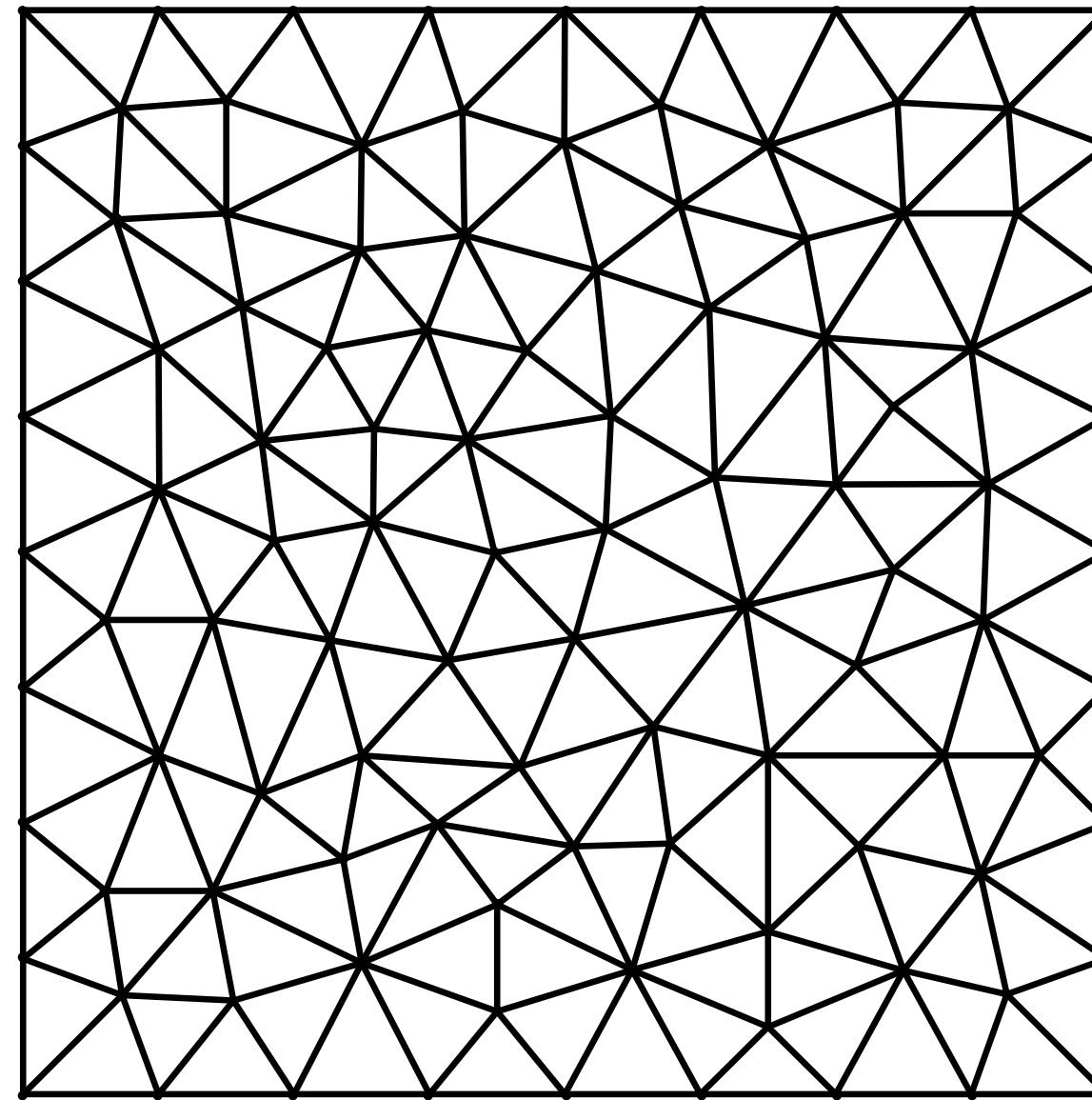
$$\nabla^2 u = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \dots$$

Spatial discretizations

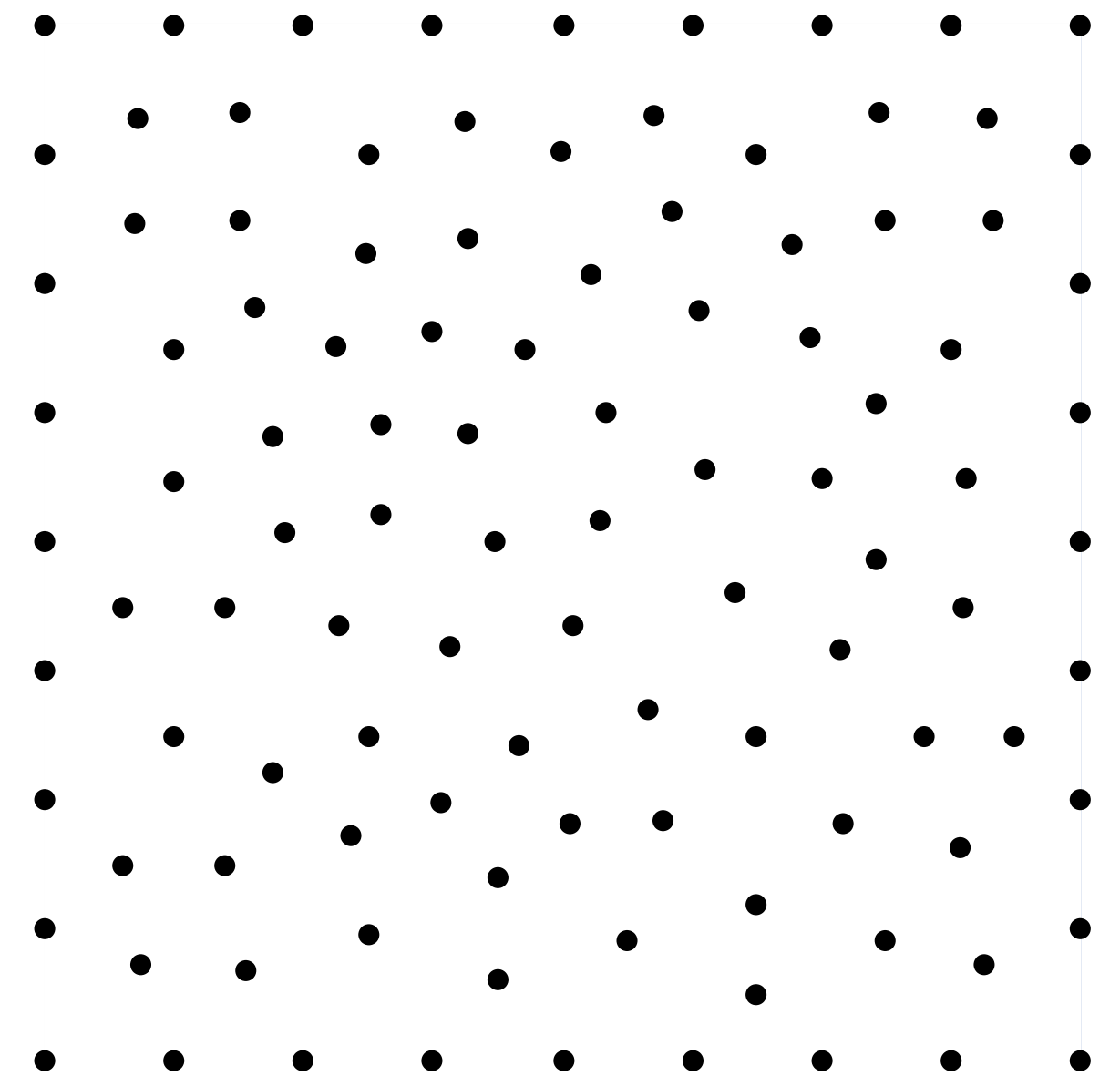
Now we have more choices of how to discretize space...



Grids



Meshes



Particles

Easier computation ←————→ More flexibility

Example: Finite differences on grids in 2D

$$\frac{\partial u}{\partial x} \approx \frac{u_{i+1,j} - u_{i,j}}{\Delta x} \quad (\text{or } \frac{u_{i,j} - u_{i-1,j}}{\Delta x} \quad \text{or } \frac{u_{i+1,j} - u_{i-1,j}}{2\Delta x} \dots)$$

$$\frac{\partial^2 u}{\partial x^2} \approx \frac{u_{i-1,j} - 2u_{i,j} + u_{i+1,j}}{\Delta x^2}$$

So the Laplacian becomes

$$\nabla^2 u = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \approx \frac{u_{i-1,j} + u_{i+1,j} + u_{i,j-1} + u_{i,j+1} - 4u_{i,j}}{\Delta x^2}$$

| | | | | |
|---|-------------|-------------|-------------|---|
| | | ⋮ | | |
| | | $u_{i,j+1}$ | | |
| ⋯ | $u_{i-1,j}$ | $u_{i,j}$ | $u_{i+1,j}$ | ⋯ |
| | | $u_{i,j-1}$ | | |
| | | ⋮ | | |

| | | | |
|---|----|---|--|
| | | 1 | |
| 1 | -4 | 1 | |
| | | 1 | |

For concreteness, here's one possible spatial discretization of the wave equation in 2D:

$$\frac{\partial^2 u}{\partial t^2} = \nabla^2 u$$

↓

$$\frac{du_{i,j}}{dt} = v_{i,j}$$

$$\frac{dv_{i,j}}{dt} = \frac{u_{i-1,j} + u_{i+1,j} + u_{i,j-1} + u_{i,j+1} - 4u_{i,j}}{\Delta x^2}$$