

<u>onstraints</u>



Last class

Backward Euler gives us a system of equations in the unknown next state (q_{n+1}, v_{n+1}):

drop the force Jacobians:

 $(\tilde{\mathbf{q}} + \Delta \mathbf{q}) = \mathbf{q}_n + (\tilde{\mathbf{v}} + \Delta \mathbf{v}) \Delta t$

What kind of time integration scheme do you get? Does it reduce to a known one?

- $\mathbf{q}^{n+1} = \mathbf{q}^n + \mathbf{v}^{n+1} \Delta t$ $v^{n+1} = v^n + M^{-1} f(q^{n+1}, v^{n+1}) \Delta t$
- Suppose you try to implement it with Newton's method starting at $\tilde{\mathbf{q}} = \mathbf{q}_n$, $\tilde{\mathbf{v}} = \mathbf{v}_n$, but you

 - $(\tilde{\mathbf{v}} + \Delta \mathbf{v}) \approx \mathbf{v}_n + \mathbf{M}^{-1} (\mathbf{f}(\tilde{\mathbf{q}}, \tilde{\mathbf{v}}) + \frac{\partial \mathbf{f}}{\partial \mathbf{r}} (\tilde{\mathbf{q}}, \tilde{\mathbf{v}}) \Delta \mathbf{q} + \frac{\partial \mathbf{f}}{\partial \mathbf{r}} (\tilde{\mathbf{q}}, \tilde{\mathbf{v}}) \Delta \mathbf{r}$

$$(\mathbf{q}_n + \Delta \mathbf{q}) = \mathbf{q}_n + (\mathbf{v}_n + \Delta \mathbf{v})$$
$$(\mathbf{v}_n + \Delta \mathbf{v}) = \mathbf{v}_n + \mathbf{M}^{-1} (\mathbf{f}(\mathbf{q}))$$

 $\mathbf{q}_{n+1} = \mathbf{q}_n + \mathbf{v}_{n+1} \Delta t$

velocity (we're using $f(q_n, v_n)$ instead of $f(q_n, v_{n+1})$)

Moral: Approximating backward Euler can still give you good behaviour.

Corollary: You can be explicit in some terms and implicit in others, for example

• use $f(q_{n+1}, v_{n+1})$ only for strong forces that cause instability

• use $f(q_n, v_n)$ for weak ones (especially if their Jacobian is hard to compute)

- **∨**) ∆t $(\mathbf{q}_n, \mathbf{v}_n) + \mathbf{0} \Delta \mathbf{q} + \mathbf{0} \Delta \mathbf{v}) \Delta t$
- $\mathbf{v}_{n+1} = \mathbf{v}_n + \mathbf{M}^{-1} \mathbf{f}(\mathbf{q}_n, \mathbf{v}_n) \Delta t$

This is basically semi-implicit Euler again! Except the acceleration term is also explicit in

Constraints

Example: How would you model a pendulum?

Make it a spring with rest length ℓ_0 , spring constant k_s , then take k_s very large?

Rule of thumb: explicit methods are only stable when $\Delta t = O(T_{\text{fast}})$, where $T_{\text{fast}} = \text{fastest timescale of dynamics in the system}$

• Period of horizontal swing: $T_{slow} \approx O(\sqrt{l_0/g})$

• Period of vertical vibration of spring: $T_{\text{fast}} \approx O(\sqrt{m/k_s})$

When k_s is very very large, T_{fast} and stable Δt become very very small!

We only care about dynamics on the scale of T_{slow} , but we're forced to take time steps on the scale of $T_{\text{fast}} \ll T_{\text{slow}}$.

In such cases, we say the problem is stiff. This happens a lot in graphics...







https://www.youtube.com/watch?v=2R9u-tjhRYA

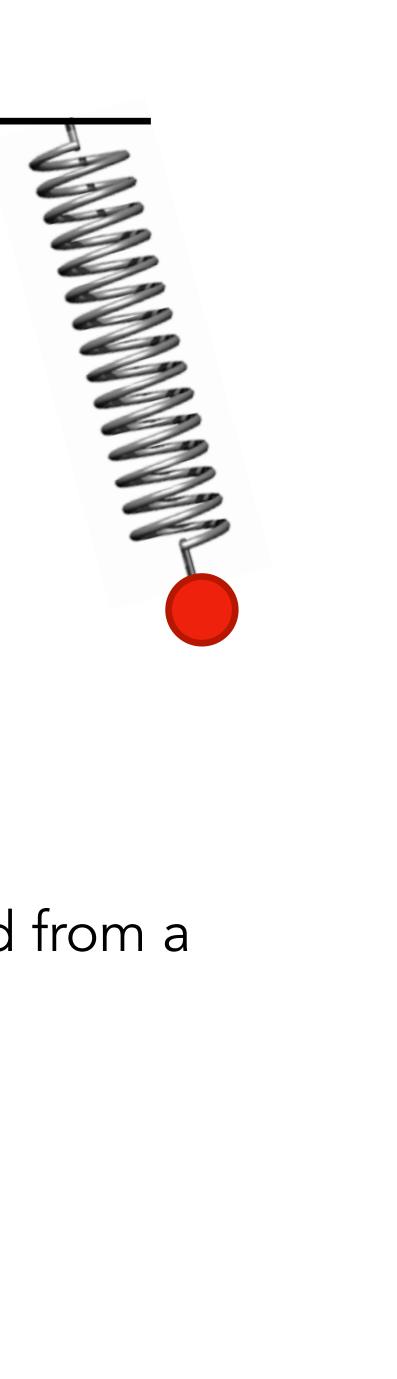
Constraints

Another general problem-solving strategy: If a parameter being very large is causing problems, make it infinity instead.

What happens to the spring when $k_s \rightarrow \infty$?

Puzzle:

- valid initial state, $\|\mathbf{x}_{ij}^{0}\| = \ell_{0}$)
- What can you say about the direction and magnitude of the spring force?



 $\mathbf{f}_{ij} = -k_s (||\mathbf{x}_{ij}|| - \ell_0) \, \mathbf{\hat{x}}_{ij}$

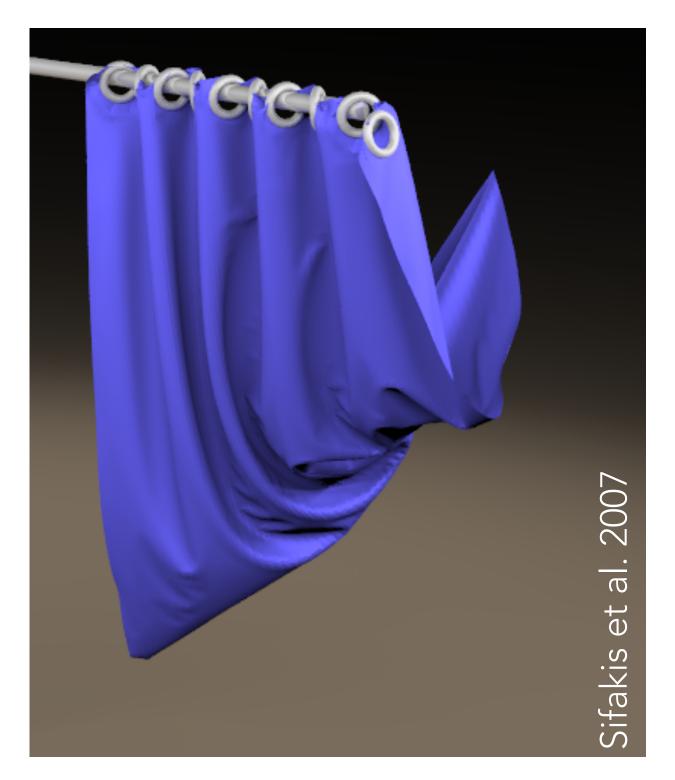
Physically, does the behaviour of this system still make sense? (At least if started from a

Original equations of motion:

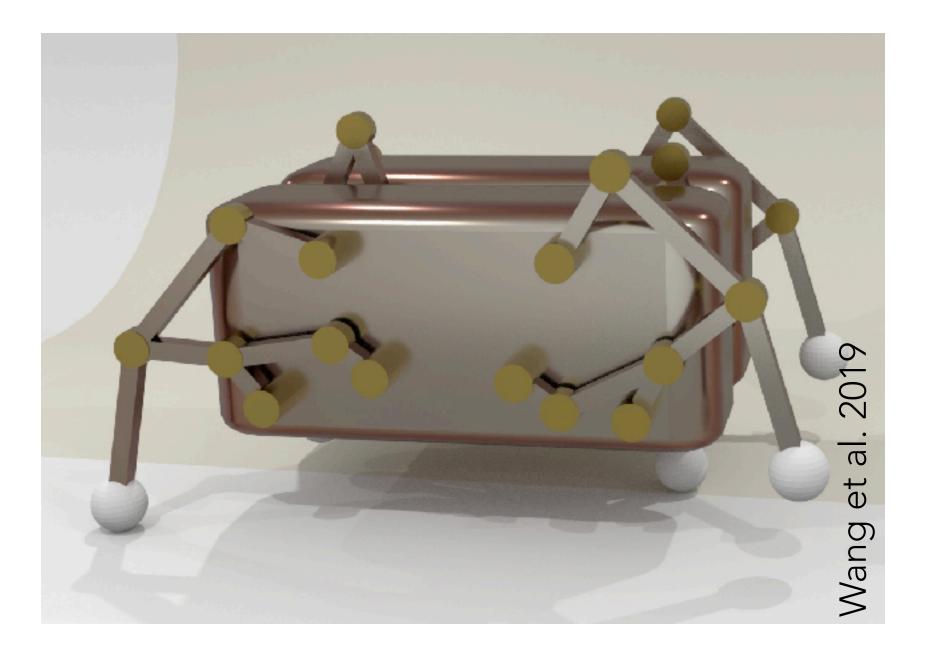
 $\ddot{\mathbf{x}} = \mathbf{g} - m^{-1} k_s (\|\mathbf{x}_{ij}\| - \ell_0) \hat{\mathbf{x}}$ Constrained equations of motion: $\ddot{\mathbf{x}} = \mathbf{g} + m^{-1} \lambda \hat{\mathbf{x}}$ $\|\mathbf{x}_{ij}\| = \ell_0$

- One new unknown: constraint force magnitude λ .
- One new equation: constraint $\|\mathbf{x}_{ij}\| = \ell_0$.

 λ is such that constraint remains satisfied over time...



Sliding on a fixed line / curve / surface



Joints between rigid parts



Inextensible cloth

In general, we may have lots of constraints on the system, each of the form

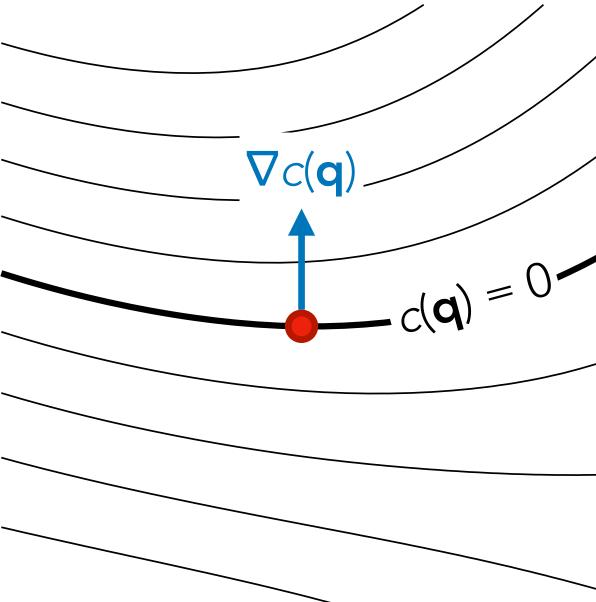
Constraint force:

Force is orthogonal to constraint surface \Rightarrow only resists moving away from constraint, not along constraint

Exercise: verify that the inextensible spring constraint from before is of this form.

- $c_i(\mathbf{q}) = 0$

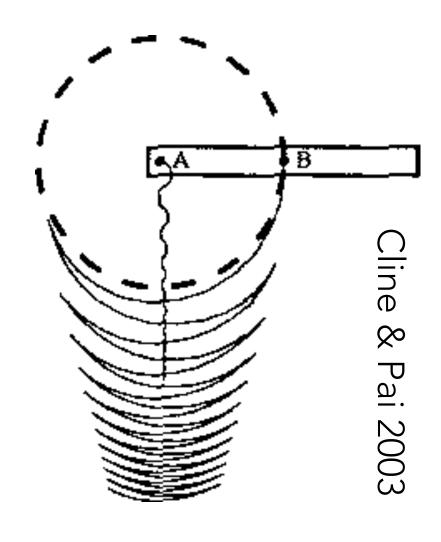
- $\mathbf{f}_{i} = \boldsymbol{\lambda}_{i} \, \nabla c_{i} \left(\mathbf{q} \right)$

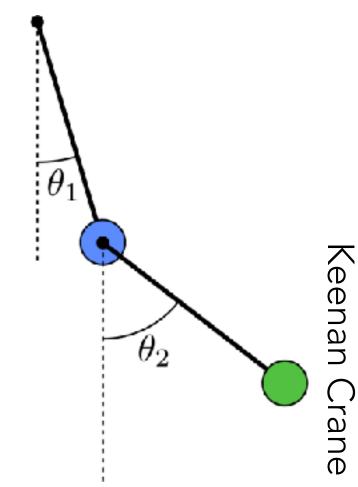


How to actually do time stepping of such a system?

- Try to estimate instantaneous λ_i at each $t_n \Rightarrow drift$
- Replace with penalty force: $\lambda_i = -k c_i(\mathbf{q}) \Rightarrow$ soft constraints
- Choose parameterization that automatically satisfies constraints ⇒ reduced coordinates
- Treat constraint forces implicitly: solve for all λ_i 's so that all $c_i(\mathbf{q}_{n+1}) = 0$

- $c_j(\mathbf{q}) = 0$ $\mathbf{f}_j = \lambda_j \nabla c_j(\mathbf{q})$
- $\ddot{\mathbf{q}} = \mathbf{M}^{-1} \left(\mathbf{f}(\mathbf{q}, \, \dot{\mathbf{q}}) + \sum_{i} \mathbf{f}_{i} \right)$





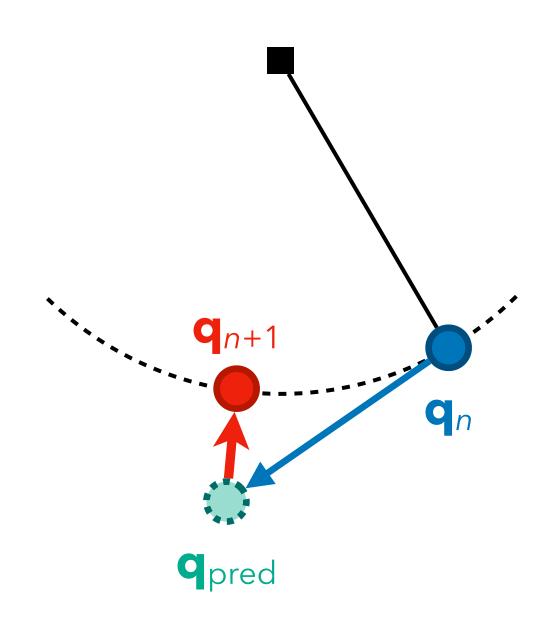
Suppose we treat the external forces explicitly and the constraint forces implicitly.

We can also eliminate \mathbf{v}_{n+1} :

Solve for \mathbf{q}_{n+1} and λ_1 , λ_2 , ... simultaneously using Newton's method

...Then update $\mathbf{v}_{n+1} = (\mathbf{q}_{n+1} - \mathbf{q}_n)/\Delta t$

- $\ddot{\mathbf{q}} = \mathbf{M}^{-1} \left(\mathbf{f}(\mathbf{q}, \dot{\mathbf{q}}) + \sum \lambda_i \nabla C_i(\mathbf{q}) \right)$
 - $c_i(\mathbf{q}) = 0$



- $\mathbf{q}_{n+1} = \mathbf{q}_{\text{pred}} + \sum \mathbf{M}^{-1} \lambda_j \nabla c_j (\mathbf{q}_{n+1}) \Delta t^2$ $c_j (\mathbf{q}_{n+1}) = 0$
- where $\mathbf{q}_{\text{pred}} = \mathbf{q}_n + \mathbf{v}_n \Delta t + \mathbf{M}^{-1} \mathbf{f}(\mathbf{q}_n, \mathbf{v}_n) \Delta t^2$.

Position-based dynamics

For real-time graphics, solving a big linear system for all λ 's is too expensive! But it's easy to solve one constraint at a time:

Example: Inextensible spring between particles *i* and *j*

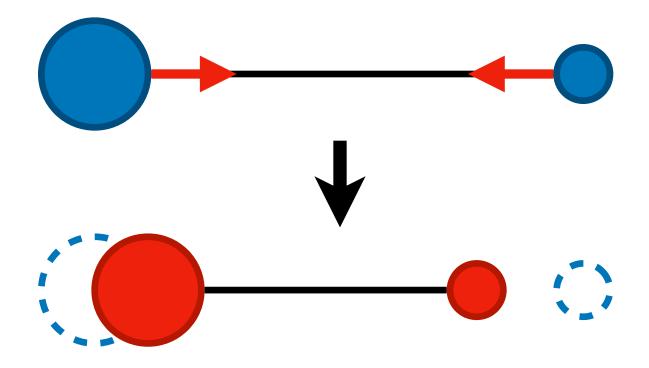
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Recall $\mathbf{q}_{n+1} = \mathbf{q}_{\text{pred}} + \sum \mathbf{M}^{-1} \lambda_j \hat{\mathbf{x}}_{ij} \Delta t^2$

Find $\Delta\lambda$ which makes updated positions satisfy $\|\mathbf{\tilde{x}}_{ij} + \Delta \mathbf{x}_{ij}\| = \ell_0$

$$\mathbf{x}_{ij} || = \ell_0$$

$$S_{ij} = \lambda \ \hat{\mathbf{x}}_{ij}$$



 $\Delta \mathbf{q}_{n+1} = \mathbf{M}^{-1} \Delta \lambda \, \hat{\mathbf{x}}_{ii} \, \Delta t^2$

In general, we have a guess of the next positions: $\tilde{\mathbf{q}}$

- 1. Applying a constraint force $\Delta \lambda_i$ changes the positions by $\Delta \mathbf{q} = \mathbf{M}^{-1} \Delta \lambda_i \nabla c_i(\mathbf{\tilde{q}}) \Delta t^2$
- 2. Solve for $\Delta \lambda_i$ so that $c_i(\tilde{\mathbf{q}} + \Delta \mathbf{q}) = 0$
- 3. Update the positions (constraint projection): $\tilde{\mathbf{q}} \leftarrow \tilde{\mathbf{q}} + \Delta \mathbf{q}$
- 4. Repeat for other constraints

Projecting one constraint makes other constraints violated!

- Loop over all constraints = 1 iteration. Have to repeat many iterations
- If not enough iterations, constraints appear soft!



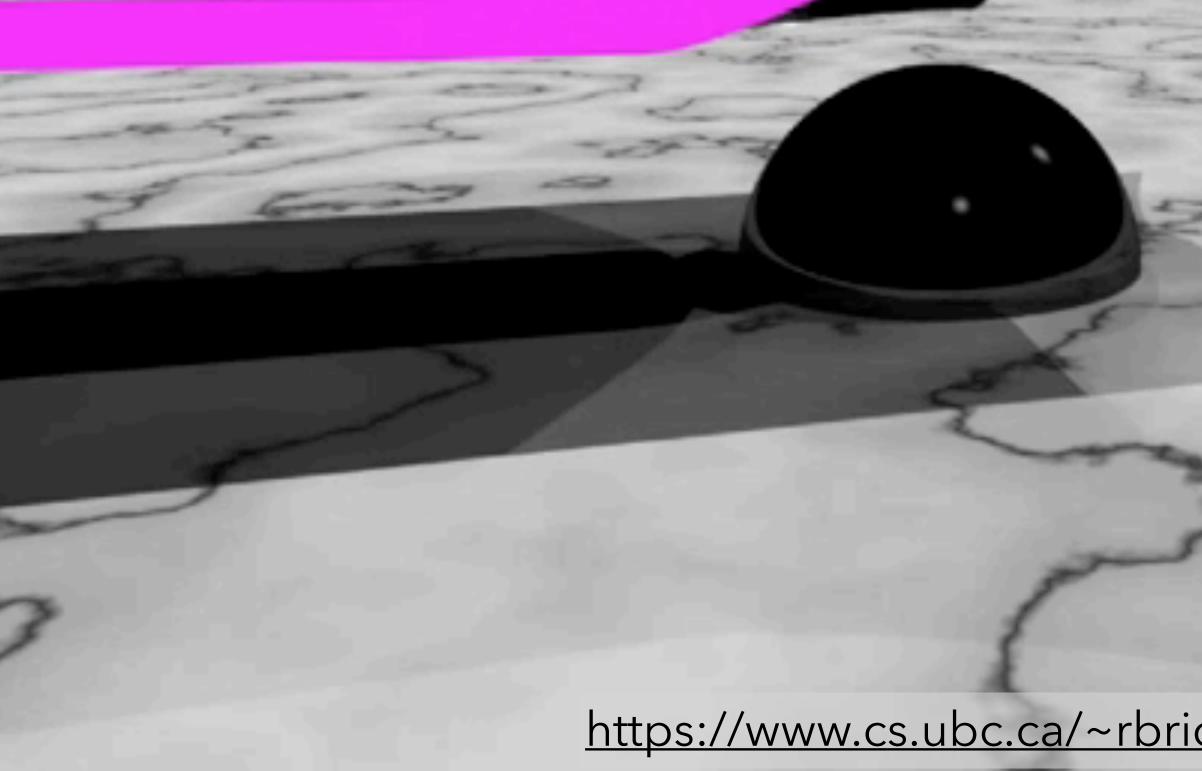
Mül ler et al. 2006



<u>https://www.youtube.com/watch?v=j5igW5-h4ZM</u>

Müller et al. 2006

Collisions



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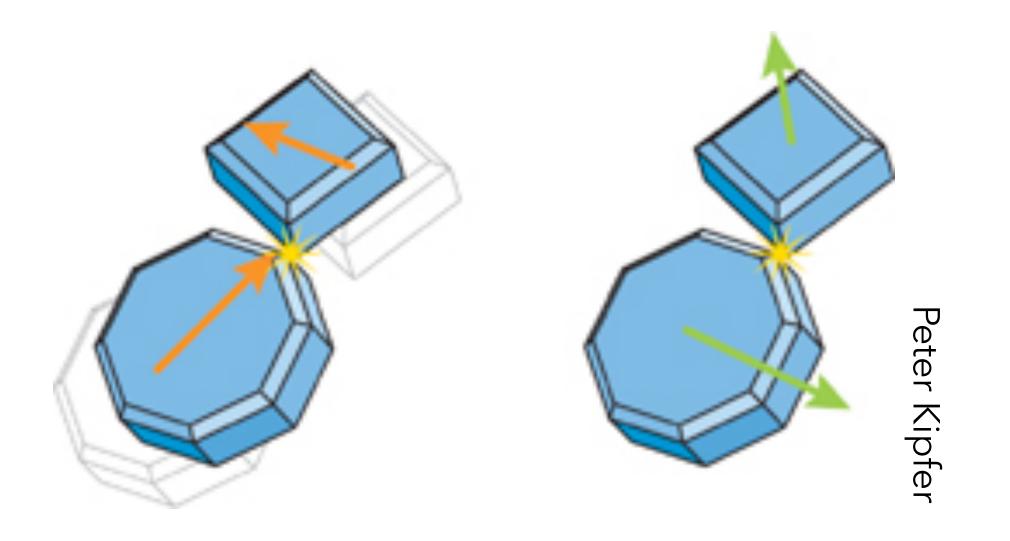
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Collision detection: find out which particles / bodies / etc. are colliding

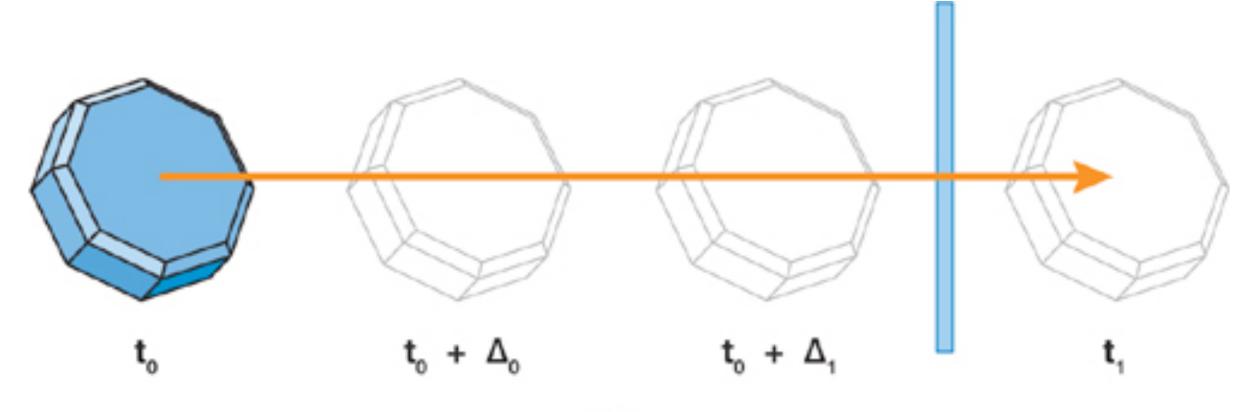
Purely a geometric problem

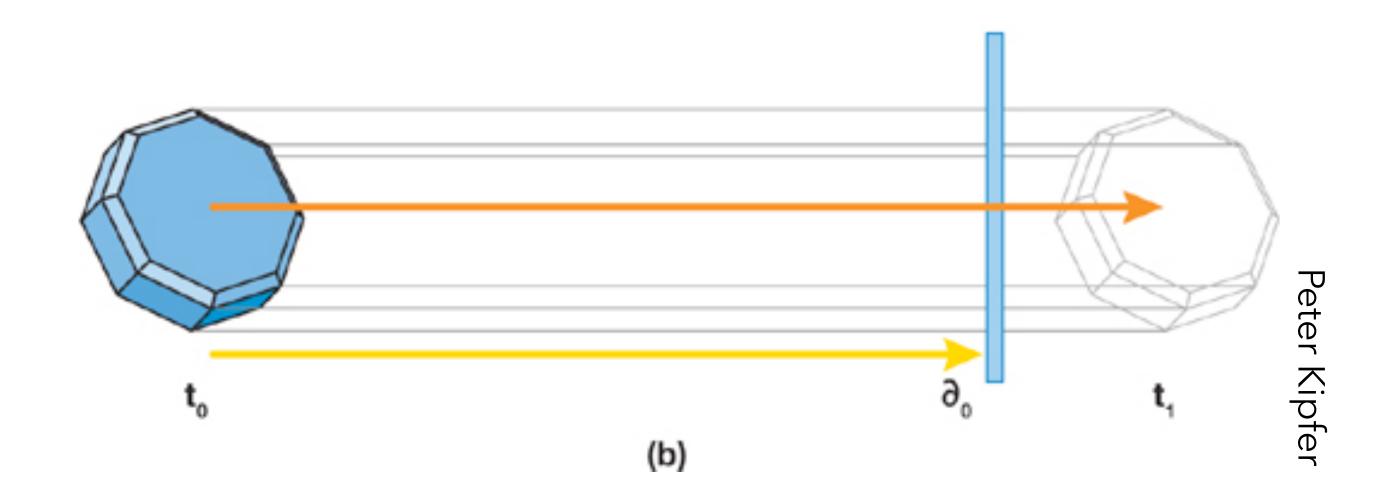


Collision response: figure out how to update their velocities / positions

Involves physics of contact forces, friction, etc.

Collision detection: discrete vs. continuous





(a)

Example: Suppose I have an infinite cylinder along the x-axis with radius R.

- I also have a particle with radius r moving to positions $\mathbf{x}_0, \mathbf{x}_1, \mathbf{x}_2, \dots$ at times t_0, t_1, t_2, \dots
- 1. How can I do discrete collision detection between the particle and the cylinder?
- 2. How can I do continuous collision detection between them?
- 3. If I model a sheet of cloth as a mass-spring system, is it enough to check that none of the particles are colliding with the cylinder?